

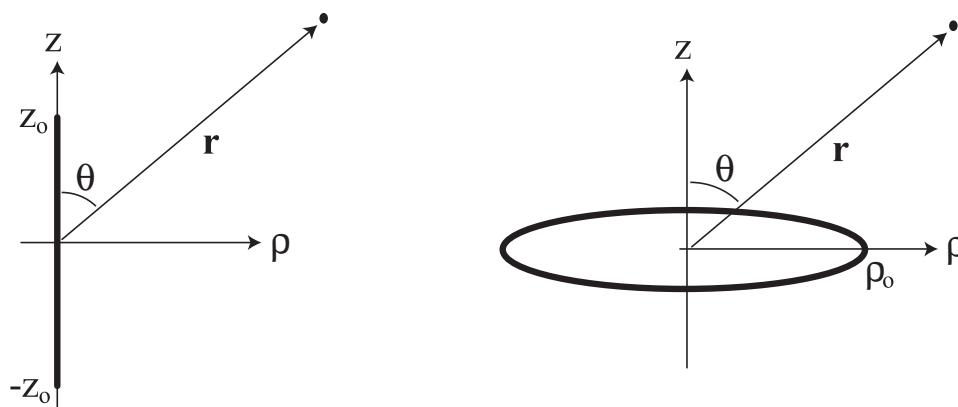
1 Gauge Transformations

The choice of $\nabla \cdot \mathbf{A}(\mathbf{r}, t)$ defines a particular gauge. With the gauge transformations we can move from one gauge to another.

- Find the gauge transformation which transforms the potentials \mathbf{A} , ϕ from the Lorentz gauge to the Coulomb gauge.

2 Line Sources and Ring Sources

Consider the time-harmonic current densities shown in the figure below. Although we cannot analytically solve for the fields we can exactly derive the farfields and the radiated power.



$$\mathbf{j}(\mathbf{r}) = I_0 \cos\left[\frac{(2n+1)\pi}{2Z_0} z\right] \delta(x)\delta(y) \mathbf{n}_z$$

$$\mathbf{j}(\mathbf{r}) = I_0 \cos(n\phi) \delta(\rho - \rho_0) \delta(z) \mathbf{n}_\phi$$

- Determine the charge densities.
- Determine expressions for the vector and scalar potentials in the Lorentz gauge and simplify as much as possible.
- At sufficiently large distances we can use the approximation

$$|\mathbf{r} - \mathbf{r}'| = \sqrt{r^2 + r'^2 - 2(\mathbf{r} \cdot \mathbf{r}')} \approx r - \mathbf{r} \cdot \mathbf{r}'/r. \quad (1)$$

Show that this approximation follows from a series expansion. The scalar Green's function can now be written as

$$G_o(\mathbf{r}, \mathbf{r}') = \frac{\exp(ik[r - \mathbf{r} \cdot \mathbf{r}'/r])}{4\pi r}. \quad (2)$$

Use this expression to calculate the vector and scalar potentials in the farfield ($r \gg r'$).

- Derive the corresponding electric and magnetic fields.

- Determine the total radiated power P for each of the two current distributions. Does it depend on the integer n ? Why?
- Choose $n=0$ for the line current and $n=1$ for the ring current. Plot the radiation patterns

$$\frac{p(\theta, \phi)}{P}, \quad (3)$$

where $p(\theta, \phi)$ is the energy flux per unit solid angle, i.e.

$$P = \int_0^\pi \int_0^{2\pi} p(\theta, \phi) \sin \theta \, d\theta \, d\phi. \quad (4)$$