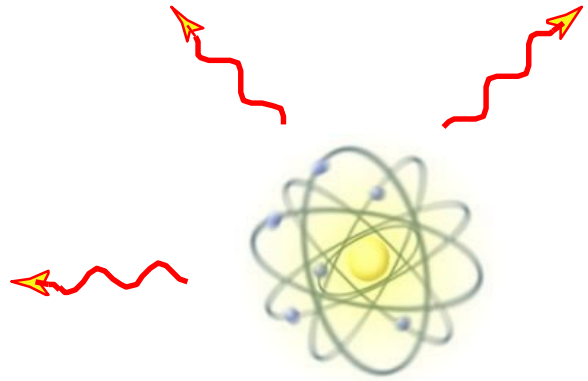
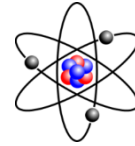


Wie entsteht Strahlung?

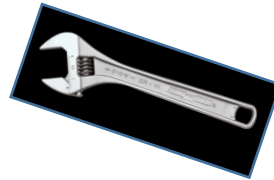


Fahrplan

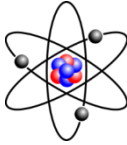
- Die inhomogene Wellengleichung



- Die Green'sche Funktion



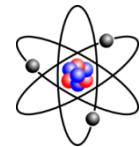
- Der elektrische Dipol – der Vater aller Strahlungsquellen



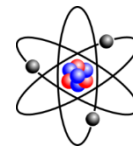
- Skalar- und Vektorpotential
 - Die Eichungen



- Die Green'sche Funktion des elektrischen Dipols



- Das Strahlungsfeld des elektrischen Dipols



Wellengleichung mit Quellterm

Inhomogene Wellengleichung:

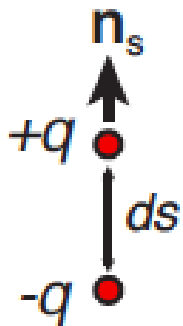
$$\nabla \times \nabla \times \mathbf{E} + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = -\mu_0 \frac{\partial}{\partial t} \left(\mathbf{j} + \frac{\partial \mathbf{P}}{\partial t} + \nabla \times \mathbf{M} \right)$$

Im monochromatischen Fall:

$$\nabla \times \nabla \times \mathbf{E}(\mathbf{r}) - k^2 \mathbf{E}(\mathbf{r}) = i\omega \mu_0 \mu(\omega) \mathbf{j}_0(\mathbf{r})$$

Für welche Inhomogenität (Quellterm) sollen wir diese (und jede andere lineare Differential-) Gleichung lösen?

Für die δ -Inhomogenität!



$$\mathbf{j}(\mathbf{r}, t) = \frac{\partial}{\partial t} \mathbf{p}(t) \delta(\mathbf{r} - \mathbf{r}_0)$$

→

$$\mathbf{j}(\mathbf{r}) = -i\omega \mathbf{p} \delta(\mathbf{r} - \mathbf{r}_0)$$

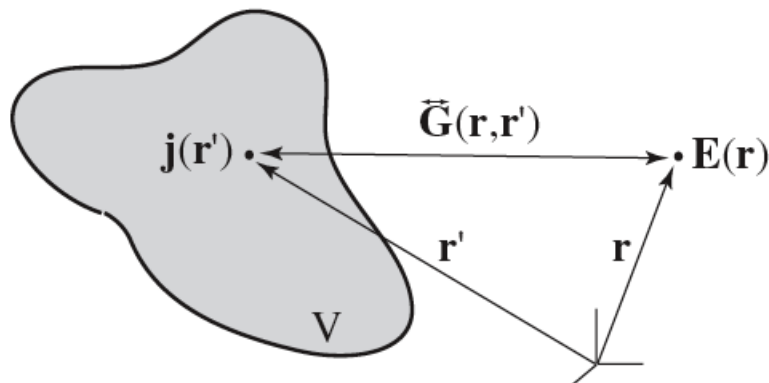
Green'sche Funktion der Wellengleichung

$$\nabla \times \nabla \times \vec{\mathbf{G}}(\mathbf{r}) - k^2 \vec{\mathbf{G}}(\mathbf{r}) = i\omega\mu_0\mu(\omega) \underbrace{(-i\omega\mathbf{p}) \perp \delta(\mathbf{r} - \mathbf{r}')}_{\mathbf{j}(\mathbf{r}) = -i\omega \mathbf{p} \delta(\mathbf{r} - \mathbf{r}_0)}$$

Mit \mathbf{G} können wir die Feldverteilung \mathbf{E} jeder beliebigen Stromverteilung berechnen!

Wie lautet \mathbf{G} ?

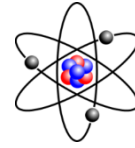
Superpositionsprinzip gilt, da Wellenoperator linear!



$$\mathbf{E}(\mathbf{r}) = \omega^2 \mu_0 \mu \vec{\mathbf{G}}(\mathbf{r}, \mathbf{r}_0) \mathbf{p}$$

Fahrplan

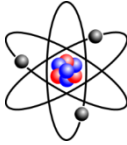
- Die inhomogene Wellengleichung



- Die Green'sche Funktion



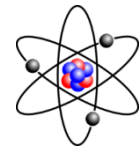
- Der elektrische Dipol – der Vater aller Strahlungsquellen



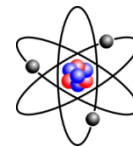
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Potentiale im monochromatischen Fall

Im monochromatischen Fall und in der Lorenzgleichung:

$$\nabla \cdot \mathbf{A} = i \frac{k^2}{\omega} \phi$$

$$\mathbf{E}(\mathbf{r}) = i\omega \mathbf{A}(\mathbf{r}) - \nabla \phi$$

$$\mathbf{B} = \nabla \times \mathbf{A}$$

$$[\nabla^2 + k^2] \mathbf{A}(\mathbf{r}) = -\mu_0 \mu \mathbf{j}_0(\mathbf{r})$$

$$[\nabla^2 + k^2] \phi(\mathbf{r}) = -\frac{1}{\epsilon_0 \epsilon} \rho_0(\mathbf{r})$$

$$[\nabla^2 + k^2] A_l(\mathbf{r}) = i\omega \mu_0 \mu \delta(\mathbf{r} - \mathbf{r}') p_l$$

$$[\nabla^2 + k^2] G_0(\mathbf{r}, \mathbf{r}') = -\delta(\mathbf{r} - \mathbf{r}')$$

Skalare Gleichung!
Helmholtz-Operator!

Überblick

Inhomogene Wellengleichung des reellen E-Feldes:

$$\nabla \times \nabla \times \mathbf{E} + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = -\mu_0 \frac{\partial}{\partial t} \left(\mathbf{j} + \frac{\partial \mathbf{P}}{\partial t} + \nabla \times \mathbf{M} \right)$$

→ Zeitharmonischer Fall:

$$\nabla \times \nabla \times \mathbf{E}(\mathbf{r}) - k^2 \mathbf{E}(\mathbf{r}) = i\omega \mu_0 \mu(\omega) \mathbf{j}_0(\mathbf{r})$$

→ Potentiale und Lorenz-Eichung

Dipol

$$[\nabla^2 + k^2] \mathbf{A}(\mathbf{r}) = -\mu_0 \mu \mathbf{j}_0(\mathbf{r})$$

$$\mathbf{j}(\mathbf{r}) = -i\omega \mathbf{p} \delta(\mathbf{r} - \mathbf{r}_0)$$

→ Green'sches Kalkül

Skalare Green'sche Funktion

$$[\nabla^2 + k^2] G_0(\mathbf{r}, \mathbf{r}') = -\delta(\mathbf{r} - \mathbf{r}')$$

$$G_0(\mathbf{r}, \mathbf{r}') = \frac{e^{ik|\mathbf{r}-\mathbf{r}'|}}{4\pi|\mathbf{r}-\mathbf{r}'|}$$

Überblick

Inhomogene Wellen

$$\nabla \times \nabla \times \mathbf{E} +$$



Zeit

$$\nabla \times \nabla \times$$



Pot

$$[\nabla^2 + k^2] A$$



Gr

$$[\nabla^2 + k^2] C$$



$$\nabla \times \mathbf{M}$$

$$j_0(\omega) \mathbf{j}_0(\mathbf{r})$$

$$-i\omega \mathbf{p} \delta(\mathbf{r} - \mathbf{r}_0)$$

$$G(\mathbf{r}, \mathbf{r}') = \frac{e^{ik|\mathbf{r}-\mathbf{r}'|}}{4\pi|\mathbf{r}-\mathbf{r}'|}$$

Überblick

Inhomogene Wellengleichung

$$\nabla \times \nabla \times \mathbf{E} + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = \mu_0 \mathbf{j} + \nabla \times \mathbf{M}$$

← Zeitharmonie

$$\nabla \times \nabla \times \mathbf{E}(\mathbf{r}, \omega) = -\mu(\omega) \mathbf{j}_0(\mathbf{r}, \omega)$$

← Potentiale

$$[\nabla^2 + k^2] \mathbf{A}(\mathbf{r}, \omega) = -\mu_0 \mathbf{j}_0(\mathbf{r}, \omega)$$

← Green'sche Funktion

$$[\nabla^2 + k^2] G_0(\mathbf{r}, \mathbf{r}') = -\delta(\mathbf{r} - \mathbf{r}') \quad G_0(\mathbf{r}, \mathbf{r}') = \frac{e^{ik|\mathbf{r}-\mathbf{r}'|}}{4\pi|\mathbf{r}-\mathbf{r}'|}$$



Und jetzt zurück zu den Feldern

Skalare Green Funktion:

$$G_0(\mathbf{r}, \mathbf{r}') = \frac{e^{ik|\mathbf{r}-\mathbf{r}'|}}{4\pi|\mathbf{r}-\mathbf{r}'|}$$

Vektorpotential:

$$\mathbf{A}(\mathbf{r}) = -i\omega\mu_0\mu G_0(\mathbf{r} - \mathbf{r}') \mathbf{p}$$

Per Definition:

$$\mathbf{E}(\mathbf{r}) = i\omega\mathbf{A}(\mathbf{r}) - \nabla\phi$$

$$\mathbf{B} = \nabla \times \mathbf{A}$$

Lorenz Eichung:

$$\nabla \cdot \mathbf{A} = i\frac{k^2}{\omega}\phi$$

Magnetisches Feld:

$$\mathbf{H}(\mathbf{r}) = -i\omega \left[\nabla \times \overset{\leftrightarrow}{\mathbf{G}}_0(\mathbf{r}, \mathbf{r}') \right] \mathbf{p}$$

Elektrisches Feld:

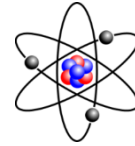
$$\mathbf{E}(\mathbf{r}) = \omega^2\mu_0\mu \overset{\leftrightarrow}{\mathbf{G}}_0(\mathbf{r}, \mathbf{r}') \mathbf{p}$$

Mit der dyadischen Green-Fkt.:

$$\overset{\leftrightarrow}{\mathbf{G}}(\mathbf{r}, \mathbf{r}') = \left[\overset{\leftrightarrow}{\mathbf{I}} + \frac{1}{k^2}\nabla\nabla \right] G_o(\mathbf{r}, \mathbf{r}')$$

Fahrplan

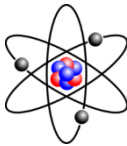
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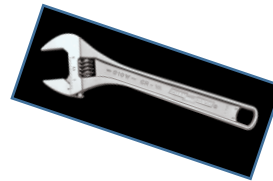
- Die Green'sche Funktion



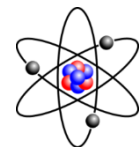
- Der elektrische Dipol – der Vater aller Strahlungsquellen



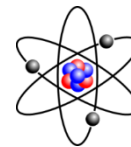
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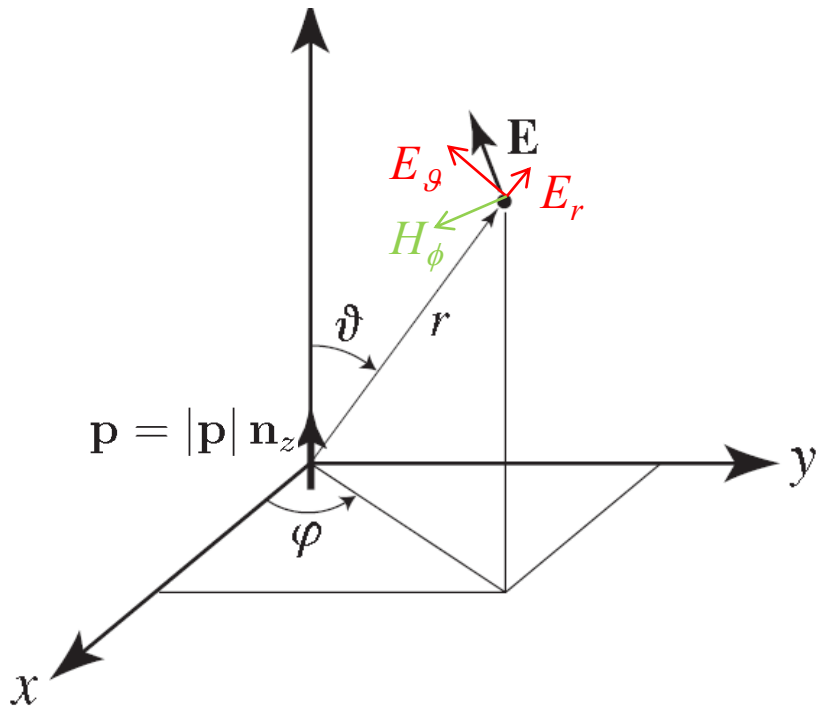
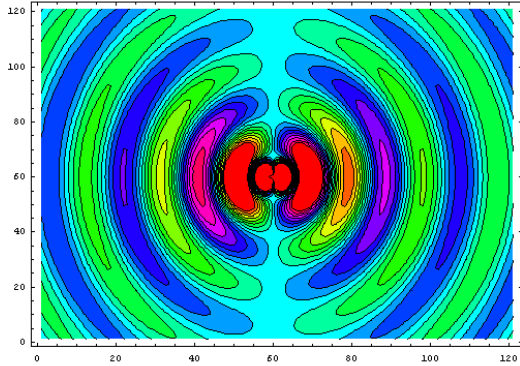
Green'sche Funktion in kartesischen Koordinaten

Mit der dyadischen Green-Fkt.: $\vec{\mathbf{G}}(\mathbf{r}, \mathbf{r}') = \left[\vec{\mathbf{I}} + \frac{1}{k^2} \nabla \nabla \right] G_o(\mathbf{r}, \mathbf{r}')$

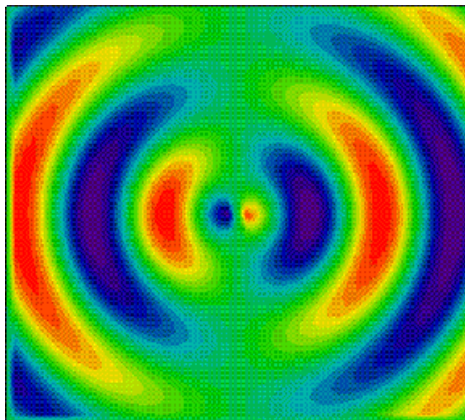
In kartesischen Koordinaten:

$$\vec{\mathbf{G}}(\mathbf{r}, \mathbf{r}_o) = \frac{\exp(ikR)}{4\pi R} \left[\left(1 + \frac{ikR - 1}{k^2 R^2} \right) \vec{\mathbf{I}} + \frac{3 - 3ikR - k^2 R^2}{k^2 R^2} \frac{\mathbf{R}\mathbf{R}}{R^2} \right]$$

Dipol-Strahlung für z-orientierten Dipol



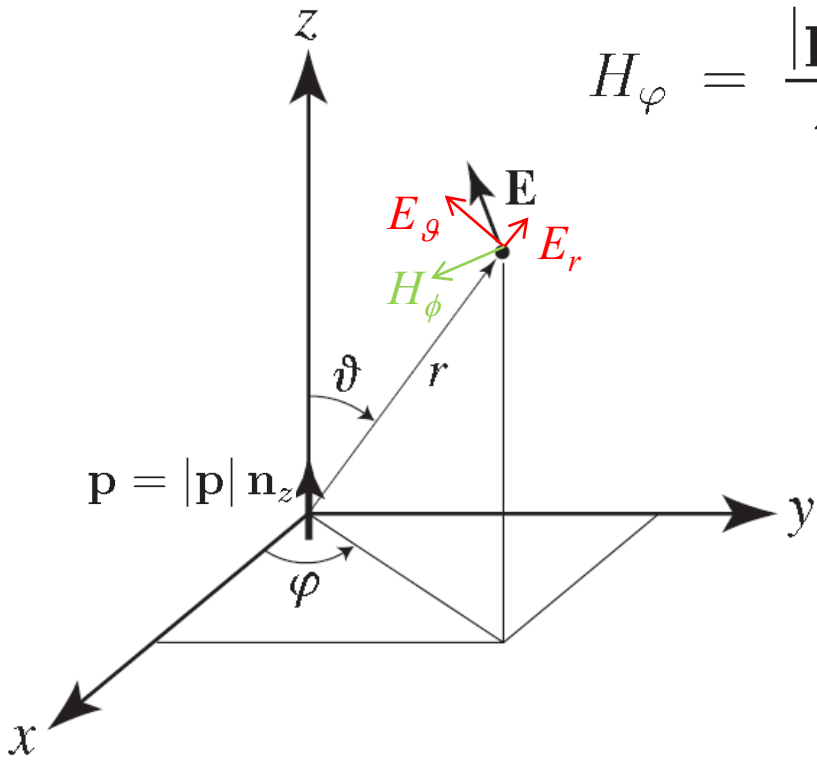
Dipol-Strahlung für z-orientierten Dipol



$$E_r = \frac{|\mathbf{p}| \cos \vartheta}{4\pi\epsilon_0\epsilon} \frac{\exp(ikr)}{r} k^2 \left[\overset{\text{NF}}{\frac{2}{k^2 r^2}} - \overset{\text{IF}}{\frac{2i}{kr}} \right],$$

$$E_\vartheta = \frac{|\mathbf{p}| \sin \vartheta}{4\pi\epsilon_0\epsilon} \frac{\exp(ikr)}{r} k^2 \left[\overset{\text{NF}}{\frac{1}{k^2 r^2}} - \overset{\text{IF}}{\frac{i}{kr}} - \overset{\text{FF}}{1} \right],$$

$$H_\varphi = \frac{|\mathbf{p}| \sin \vartheta}{4\pi\epsilon_0\epsilon} \frac{\exp(ikr)}{r} k^2 \left[\overset{\text{IF}}{-\frac{i}{kr}} - \overset{\text{FF}}{1} \right] \sqrt{\frac{\epsilon_0\epsilon}{\mu_0\mu}}$$



- NB:
- Magnetfeld hat keinen Nahfeldterm
 - Fernfelder sind transversal
 - Intermediate Field ist 90° ausser Phase mit Nah- und Fernfeld